



## Elastic properties of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ at high temperature superconductor

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### Abstract

Elastic properties at high temperature Superconductors depends on the second order (SOE) and third order (TOE) expressions. The superconductor of  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  having tetragonal structure have been derived using the finite deformation theory and variation with Einstein temperature ( $\theta_E$ ), Debye temperature ( $\theta_D$ ) for  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  at high temperature superconductor. The derived expression have been used to compute the values of SOE and TOE constant of  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  material are being presented for the first time.

**Keywords:** elastic properties, second order (SOE), third order (TOE)

### Introduction

Depending upon the response of a superconducting material to and applied magnetic field, superconductors are normally classified into two categories, type I<sup>st</sup> in which the expulsion of the magnetic lines of force is complete below the applied critical field and type II<sup>nd</sup> in which the magnetic field penetration is complete till the upper critical field destroys superconducting [1]. The longitudinal sound velocity in the Barium doped compound with the composition  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  was measured by Fossheim *et al.* [2] using the pulse method at 24 MHz These authors also observed a sudden drop in the velocity starting from 200 to 240K. This was accompanied by a sharp rise in attenuation which exhibited a broad peak at 100 K. Copper becomes strictly square planer coordinator with Cu-O bonds of about 1.97 Å. They are no epical oxygen the La atom bond to 8 oxygen compare to 9-oxygen for larger Lanthanum in the T-structure.

Our study of elastic constant for the tetragonal phase of  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  will provide some information about the structural phase transition and doping. Lanthanum Copper Oxide  $\text{La}_2\text{CuO}_4$  undergoes a structure transformation from a tetragonal to an orthorhombic structure at 530 K. The Strontium and Barium doped superconducting compounds are tetragonal at room temperature. The phase diagram of tetragonal to orthorhombic transition in the Strontium doped compound has been investigator by Moret *et al.* [3-7].

The superconductivity on these compounds is highly sensitive to structure perturbations induced by hydrostatic pressure owing to their anisotropic structure [8-9]. We attribute the deviation from Poission behavior and the C–lattice contraction to tensile strain induced oxygen vacancy formation owing to the reduced oxygen activation energy [10-11]. This is to be compared with a Debye temperature of 292 K determined via specific heat for a superconducting single crystal elastic data are consistent with specific heat data obtained on ceramics [12-14], where values nearer to 360 K are observed. Thus the discrepancy, though suggestive of porosity and ceramic structure effects is certainly real because it is hard to imagine

how this compound could a  $V_l$  of 7.3 km/s and a  $V_s$  of 4.1 km/s, the values predicted by the specific heat data if we use  $\sigma = 0.26$ . A highly anisotropy structure could, however, limit the specific has contribution from the acoustic phonons. The super conducting phase transition generates a discontinuity in compression elastic constants which can be obtained from a thermodynamic treatment. Observation of such effects are relevant to the interpretation of normal state properties because the discontinuity.  $\Delta V_l$  is the bulk sound velocity of an isotropic superconductor given by-  $\Delta V_l = \frac{V_l \Delta c}{2T_c} \left[ \frac{dT_c}{dP} \right]^2$ , where  $\Delta c$  is the discontinuity in specific heat at  $T_c$  and  $V_l$  is the bulk sound speed.

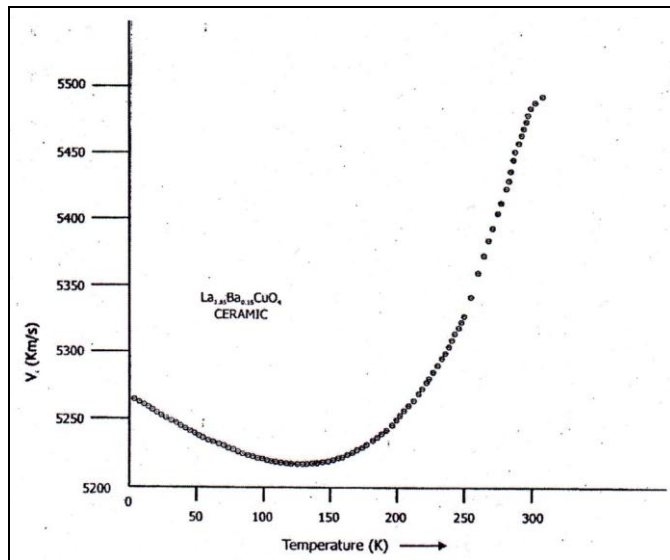
### Calculation and Results

**Table 1:** Value of Einstein Temperature ( $\theta_E$ ) and Debye Temperature ( $\theta_D$ ) for various high temperature superconductors calculated from specific heat data

Compounds	T(K)	$\frac{C}{3R}$	( $\theta_E$ )(K)	( $\theta_D$ )(K)
$\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$	10	0.020	80.9	157.3
	15	0.060	98.8	162.8
	20	0.136	107.2	161.0
	25	0.261	107.0	151.3
	30	0.441	97.8	163.0
	35	0.646	82.0	108.5
	40	0.750	75.4	96.1

**Table 2:** Second Order Elastic Constant (in GPa) of  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$

$C_{ij}$	Present
$C_{11}$	239.89
$C_{12}$	84.01
$C_{13}$	55.01
$C_{33}$	147.06
$C_{44}$	53.97
$C_{66}$	80.22



**Fig 1:**  $V_L$  for single crystal  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  ceramic showing the decreases in sound velocity on cooling below 300K.

### Result and Discussion

The second order elastic constants thus obtained for  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  ( $x=0.10$  to  $0.20$ ) are given in above table. A comparison of second order elastic constants calculated using the present method and the measured value for  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  superconductors are also given in table. The agreement of second order elastic constants obtained by present method with the some reported experimental value is good.

The third order elastic constants obtained by us are listed. The third order elastic constants are compared with those from <sup>[15]</sup>, where in the values are given as combinations of third order elastic constants. For  $C_{11}+6C_{12}+2C_{23}$  and  $C_{44}+2C_{66}$  they obtained  $-301 \times 10^{11} \text{ Pa}$  and  $-6 \times 10^{11} \text{ Pa}$  respectively. The corresponding values obtained by us are  $-90.87 \times 10^{11} \text{ Pa}$  and  $-24.78 \times 10^{11} \text{ Pa}$  respectively. The CuO planes do not suffer relative displacement when the lattice is strained. All the third order elastic constants of  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  ( $x=0.13$  to  $0.20$ ) are negative.

The change in elastic properties is undoubtedly related to the tetragonal to orthorhombic transition at  $T_0$ . In fact, coupling between the order parameter, and strains can account for this softening quite naturally. However, the acoustic modes remain soft for an extended temperature range below  $T_0$ . One major objects of the present study has also been to analysis the elastic behavior on other ceramic materials to find out whether a small bulk modulus and larger pressure dependence are general characteristics of mixed oxides based on copper.

### References

1. Lechevet J, Neighbor JE, Shiffman CA. J. Low Temp. Phys. 1977; 27:407.
2. Fossheim K, Laegried T, Sandvold E, Vassenden F, Muller KA, Bednorz JG. Phys. Rev. Letters. 1987; 58:1143.
3. Moret R, Pouget JP, Collin G, Europhys. Lrtt. 1987; 4:365.
4. Muller V, Maurer D, Roth CH, Hucho C, Winau D, De Groot K, *et al.* Physica. 1988; C153-155:280.
5. Srinivasan R, Ramachandran V, Seshadri AT, Ananda Ramadass G, Pramana. J.Phys. 1987; 29:L603.
6. Hardy RJ, Binek C. Thermodynamics and statistical Mechanics: An Integrated Approach (Wiley 2014).
7. Binek C. Refrigeration through voltage-controlled entropy change. U.S. patent 9366460B2, 2016.
8. Bud'ko SL, Guimpel J, Nakamura O, Maple MB, Schuller IK, Phys.Rev.B46. 1992; 1257.
9. Gugenberger F, Meingast C, Roth G, Grube K, Breit V, Wühl H, Phys. Rev. B49. 1994; 13137.
10. Mayer TL, Jiang L, Lee J, Yoon M, Freeland JW, Jang JH, *et al.* Lee e-print. arXiv: 1508.06971, 2015.
11. Aschauer U, Pfenninger R, Selbeck SM, Grand T, Spaldin NA. Phys. Rev. B88, 054111, 2013.
12. Schirber JE, Morosin B, Merrill RM, *et al.* Physica, C152, 1988.
13. Birgeneau RJ, Chen CY, Gabbe DR, *et al.* Phys. Rev. Lett. 1987; 59:1329.
14. Yu RC, Naughton MJ, Yan X, *et al.* Phys. Rev., B37. 1988, 7963.
15. Jayachandran KP, Menon CS. Indian Journal of Pure & Applied Physics. 1998; 36:188.