



## Unification of different numerical methods for the solution of linear fractional differential equation

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### Abstract

In this paper, an attempt has been made to obtain the solution of a linear ordinary fractional differential equation by Generalized Differential Transform Method (GDTM), Homotopy Analysis Method (HAM), Adomian Decomposition Method (ADM) and Power Series Method (PSM). Differential equations of fractional order, as generalization of classical differential equations of integer order, are increasingly used to model problems in fluid flow, finance and other areas of applications. Fractional derivatives are considered in the Caputo sense. Using the present methods we can solve many linear and nonlinear fractional differential equations. These methods perform extremely well in terms of accuracy, efficiency and simplicity and provide an analytical solution in the form of an infinite power series with easily computable components.

**Keywords:** fractional linear differential equation, generalized differential transform method, homotopy analysis method, adomian decomposition method, power series method, caputo fractional derivative, analytic solution

### Introduction

In recent years fractional derivatives have found to be valuable tools in the modeling of any physical phenomena in various fields of science and engineering, finance and hydrology (see Podlubny, 1999) [38]. It has been observed that fractional derivatives provide an excellent instrument for the description of memory and hereditary properties of various materials and processes. Zhou (1986) [53] proposed the differential transform method (DTM) to solve linear and nonlinear initial value problems in electric circuit analysis. Differential transform method (DTM) constructs an analytical solution in the form of a polynomial and different from the traditional higher order Taylor series method. This method has been used for solving various types of equations by many authors (see Chen and Ho, 1999; Jang and Chen, 2001; Ayaz, 2004; Arikoglu and Ozkol, 2006; Arikoglu and Ozkol, 2007; Momani *et al.*, 2008; Tari *et al.*, 2009; Kangalgil and Ayaz, 2009; Abazari and Borhanifar, 2010; Nazari and Shahmorad, 2010; Borhanifar and Abazari, 2010; Soltanalizadeh and Zarebnia, 2011; Borhanifar and Abazari, 2011) [17, 24, 10, 8, 9, 33, 36, 37, 43, 25, 1, 35, 14, 40, 15]. For solving two-dimensional linear and nonlinear partial differential equations of fractional order DTM is further developed as generalized differential transform method (GDTM) (see Momani *et al.*, 2007; Odibat and Momani, 2008; Odibat *et al.*, 2008) [34, 33, 36, 37]. Recently, Ertirka and Momanib (2010) [19] applied generalized differential transform method to solve fractional integro-differential equations. Garg *et al.* (2011) [20] implemented generalized differential transform method (GDTM) to derive the solution of space-time fractional telegraph equation. Bansal and Jain (2015) [12] applied generalized differential transform method to solve fractional order Riccati differential equation. A new analytical approach that can be applied to solve nonlinear fractional differential equations is to employ the homotopy analysis method (HAM) (see Liao, 1992, 1995, 1997, 2004, 2009) [26, 27, 28, 29, 30, 52]. Liao (1992) [26] proposed the homotopy analysis method (HAM), is a powerful method to solve the nonlinear problems. Homotopy analysis method (HAM) provides a simple way to adjust and control the convergence of the solution series by introducing a parameter  $\hbar$ , known as convergence-control parameter. This method has also been successfully applied to solve many types of nonlinear problems by several authors (see Song and Zhang, 2007; Abdulaziz *et al.*, 2008; Alomari *et al.*, 2008; Xu and Cang, 2008; Hashim *et al.*, 2009; Song and Zhang, 2009; Xu *et al.*, 2009; Jafari and Seifi, 2009; Jafari and Seifi, 2009) [41, 2, 7, 51, 21, 42, 52, 22, 23]. Adomian (1988, 1994) [3, 4] introduced the adomian decomposition method (ADM), possesses a great potential in solving different kinds of functional equations. Both linear and nonlinear equations and systems of such types are all amenable to the method. This method has the advantage of dealing directly with the problem in case of nonlinear differential equations and nonlinear partial differential equations, and equations are solved without transforming them to more simple ones. The adomian decomposition method avoids linearization, perturbation, discretization, or any unrealistic assumptions. In case of nonlinear equations, the nonlinearity term is replaced by a series of what are called Adomian polynomials and the evaluation of these polynomials is essential, as they contribute to the solution's series components. Thus, relations and algorithms have been deduced and continuously improved to obtain such polynomials in an easy way (see Wazwaz, 2000; Biazar *et al.*, 2003; Babolian and Javadi, 2004; Zhu *et al.*, 2005; Pourdarvish, 2006) [45, 46, 13, 11, 54, 39].

**2. Formation of the problem**

Consider the linear ordinary fractional differential equation:

$$\frac{dy(x)}{dx} + \frac{d^{1/2}y(x)}{dx^{1/2}} + y(x) = x \quad ; \quad x \in [0,1] \tag{1}$$

subject to the initial conditions  $y(0) = 1$  and  $y'(0) = -1$  (2)

where  $\frac{d^{1/2}}{dx^{1/2}}$  is the fractional differential operator (Caputo derivative) of order  $1/2$ .

**3. Mathematical Preliminaries on Fractional Calculus**

In the present analysis we introduce the following definitions (see Miller and Ross, 1993; Das, 2008) <sup>[32, 18]</sup>.

**3.1 Definition:** Let  $\alpha \in R^+$ . On the usual Lebesgue space  $L(a,b)$  integral operator  $I^\alpha$  defined by

$$I^\alpha f(x) = \frac{d^{-\alpha} f(x)}{dx^{-\alpha}} = \frac{1}{\Gamma(\alpha)} \int_0^x (x-t)^{\alpha-1} f(t) dt \quad \text{and} \quad I^0 f(x) = f(x)$$

is called Riemann-Liouville fractional integral operator of order  $\alpha \geq 0$  and  $a \leq x < b$

It has the following properties:

$$I^\alpha f(x) \text{ exists for any } x \in [a,b]$$

$$I^\alpha I^\beta f(x) = I^{\alpha+\beta} f(x)$$

$$I^\alpha I^\beta f(x) = I^\beta I^\alpha f(x)$$

$$I^\alpha x^\gamma = \frac{\Gamma(\gamma+1)}{\Gamma(\alpha+\gamma+1)} x^{\alpha+\gamma}$$

Where  $f(x) \in L[a,b]$ ,  $\alpha, \beta \geq 0$ ,  $\gamma > -1$

**3.2 Definition:** The Riemann-Liouville definition of fractional order derivative is

$${}^{RL}D_x^\alpha f(x) = \frac{d^n}{dx^n} I_x^{n-\alpha} f(x) = \frac{1}{\Gamma(n-\alpha)} \frac{d^n}{dx^n} \int_0^x (x-t)^{n-\alpha-1} f(t) dt,$$

where  $n$  is an integer that satisfies  $n-1 < \alpha < n$ .

**3.3 Definition:** A modified fractional differential operator  ${}^cD_x^\alpha$  proposed by Caputo is given by

$${}^cD_x^\alpha f(x) = I_x^{n-\alpha} \frac{d^n}{dx^n} f(x) = \frac{1}{\Gamma(n-\alpha)} \int_0^x (x-t)^{n-\alpha-1} f^{(n)}(t) dt,$$

Where  $\alpha (\alpha \in R^+)$  is the order of operation and  $n$  is an integer that satisfies  $n-1 < \alpha < n$ .

It has the following two basic properties ( see Caputo, 1967):

$$\text{If } f \in L_\infty(a, b) \text{ or } f \in C[a, b] \text{ and } \alpha > 0 \text{ then } {}^c D_x^\alpha I_x^\alpha f(x) = f(x).$$

$$\text{If } f \in C^n[a, b] \text{ and if } \alpha > 0 \text{ then}$$

$${}_0 I_x^\alpha {}^c D_x^\alpha f(x) = f(x) - \sum_{k=0}^{n-1} \frac{f^{(k)}(0^+)}{k!} x^k, \quad n-1 < \alpha < n.$$

**3.4 Definition:** For  $m$  being the smallest integer that exceeds  $\alpha$ , the Caputo time-fractional derivative operator of order  $\alpha > 0$ , is defined as (see Podlubny, 1999) [38]

$$D_t^\alpha u(x, t) = \frac{\partial^\alpha u(x, t)}{\partial t^\alpha} = \begin{cases} \frac{\partial^m u(x, \xi)}{\partial \xi^m} & ; \alpha = m \in N \\ \frac{1}{\Gamma(m-\alpha)} \int_0^t (t-\xi)^{m-\alpha-1} \frac{\partial^m u(x, \xi)}{\partial \xi^m} d\xi & ; m-1 \leq \alpha < m \end{cases}$$

**Relation between Caputo derivative and Riemann-Liouville derivative:**

$${}^c D_x^\alpha f(x) = {}^{RL} D_x^\alpha f(x) - \sum_{k=0}^{m-1} \frac{f^{(k)}(0^+)}{\Gamma(k-\alpha+1)} x^{k-\alpha}, \quad m-1 < \alpha < m$$

Integrating by parts, we get the following formulae as given by ( see Almeida and Torres, 2011)

$$\int_a^b g(x) {}^c D_x^\alpha f(x) dx = \int_a^b f(x) {}^{RL} D_b^\alpha g(x) dx + \sum_{j=0}^{n-1} \left[ {}^{RL} D_b^{\alpha+j-n} g(x) {}^{RL} D_b^{n-j-1} f(x) \right]_a^b$$

$$\text{For } n=1, \int_a^b g(x) {}^c D_x^\alpha f(x) dx = \int_a^b f(x) {}^{RL} D_b^\alpha g(x) dx + \left[ {}_x I_b^{1-\alpha} g(x) \cdot f(x) \right]_a^b$$

**4. Generalized one dimensional differential transform method:**

Generalized differential transform of a function  $y(x)$  in one variable is denoted by  $Y_\alpha(k)$  and defined as follows (see Momani *et al.*, 2007; Odibat and Momani, 2008; Odibat *et al.*, 2008; Ertiirka and Momanib, 2010; Garg *et al.*, 2011; Bansal and Jain, 2015) [34, 33, 36, 37, 20, 12].

$$Y_\alpha(k) = \frac{1}{\Gamma(\alpha k + 1)} \left[ \left( D_{x_0}^\alpha \right)^k y(x) \right]_{x=x_0} \tag{3}$$

where  $\alpha \in (0, 1]$  and  $\left( D_{x_0}^\alpha \right)^k = D_{x_0}^\alpha, D_{x_0}^\alpha, \dots, D_{x_0}^\alpha$  ( $k$  - times).

and the inverse generalized differential transform of  $Y_\alpha(k)$  is given by

$$y(x) = \sum_{k=0}^{\infty} Y_\alpha(k) (x - x_0)^{\alpha k} \tag{4}$$

It has the following properties:

$$\text{If } u(x) = v(x) \pm w(x) \text{ then } U_\alpha(k) = V_\alpha(k) \pm W_\alpha(k)$$

$$\text{If } u(x) = av(x); a \in R \text{ then } U_\alpha(k) = aV_\alpha(k)$$

$$\text{If } U(x) = v(x)w(x) \text{ then } U_\alpha(x) = \sum_{r=0}^k V_\alpha(r)W_\alpha(k-r)$$

$$\text{If } u(x) = (x-x_0)^{n\alpha} \text{ then } U_\alpha(k) = \delta(k-n)$$

$$\text{If } u(x) = D_{x_0}^\alpha v(x); 0 < \alpha \leq 1 \text{ then } U_\alpha(k) = \frac{\Gamma(\alpha(k+1)+1)}{\Gamma(\alpha k+1)} V_\alpha(k+1)$$

If  $u(x) = x^\lambda f(x)$  where  $\lambda > -1$ ,  $f(x)$  has the generalized Taylor series expansion  $f(x) = \sum_{n=0}^{\infty} a_n (x-x_0)^{n\alpha}$  and

$\beta < \lambda + 1$  and  $\alpha$  is arbitrary or

$\beta \geq \lambda + 1$ ,  $\alpha$  arbitrary and  $a_n = 0$  for  $n = 0, 1, 2, \dots, m-1$ , where  $m-1 < \beta \leq m$ . Then (3) becomes

$$U_\alpha(k) = \frac{1}{\Gamma(\alpha k+1)} [D_{x_0}^{\alpha k} u(x)]_{x_0}$$

If  $u(x) = D_{x_0}^\gamma f(x)$ ,  $m-1 < \gamma \leq m$  and the function  $f(x)$  satisfies the conditions given in (VI) then

$$U_\alpha(k) = \frac{\Gamma(\alpha k + \gamma + 1)}{\Gamma(\alpha k + 1)} F_\alpha\left(k + \frac{\gamma}{\alpha}\right)$$

where  $U_\alpha(k)$ ,  $V_\alpha(k)$ ,  $W_\alpha(k)$  and  $F_\alpha(k)$  are the differential transformations of the functions  $u(x)$ ,  $v(x)$ ,  $w(x)$  and  $f(x)$  respectively and

$$\delta(k-n) = \begin{cases} 1 & ; k = n \\ 0 & ; k \neq n \end{cases}$$

Applying generalized differential transform (3) with  $x_0 = 0$  on (1) and (2) we obtain

$$Y_{\frac{1}{2}}(k+2) = \frac{\Gamma\left(\frac{k}{2}+1\right)}{\Gamma\left(\frac{k}{2}+2\right)} \left[ \delta(k-2) - Y_{\frac{1}{2}}(k) - \frac{\Gamma\left(\frac{k}{2}+\frac{3}{2}\right)}{\Gamma\left(\frac{k}{2}+1\right)} Y_{\frac{1}{2}}(k+1) \right] \tag{5}$$

and  $Y_{\frac{1}{2}}(0)=1, Y_{\frac{1}{2}}(1)=0, Y_{\frac{1}{2}}(2)=-1$  (6)

Now utilizing the recurrence relation (5) and the initial conditions (6), we obtain after a little simplification the following values of  $Y_{\frac{1}{2}}(k)$  for  $k=0,1,2,\dots$

$$\begin{aligned}
 Y_{\frac{1}{2}}(3) &= \frac{1}{\Gamma(5/2)}, & Y_{\frac{1}{2}}(4) &= \frac{1}{\Gamma(3)}, & Y_{\frac{1}{2}}(5) &= -\frac{2}{\Gamma(7/2)}, & Y_{\frac{1}{2}}(6) &= \frac{1}{\Gamma(4)}, & Y_{\frac{1}{2}}(7) &= \frac{1}{\Gamma(9/2)}, \\
 Y_{\frac{1}{2}}(8) &= -\frac{2}{\Gamma(5)}, & Y_{\frac{1}{2}}(9) &= \frac{1}{\Gamma(11/2)}, & Y_{\frac{1}{2}}(10) &= \frac{1}{\Gamma(6)}, & Y_{\frac{1}{2}}(11) &= -\frac{2}{\Gamma(13/2)}, & Y_{\frac{1}{2}}(12) &= \frac{1}{\Gamma(7)}, \\
 Y_{\frac{1}{2}}(13) &= \frac{1}{\Gamma(15/2)}, & Y_{\frac{1}{2}}(14) &= -\frac{2}{\Gamma(8)}, & Y_{\frac{1}{2}}(15) &= \frac{1}{\Gamma(17/2)}, & Y_{\frac{1}{2}}(16) &= \frac{1}{\Gamma(9)}, & Y_{\frac{1}{2}}(17) &= -\frac{2}{\Gamma(19/2)}
 \end{aligned}$$

and so on Now, from (4), we have

$$y(x) = \sum_{k=0}^{\infty} Y_{\frac{1}{2}}(k) x^{k/2} \tag{7}$$

Using the above values of  $Y_{\frac{1}{2}}(k)$ ;  $k=0,1,2,\dots$  in (4) the solution of (1) is obtained as

$$\begin{aligned}
 y(x) &= 1-x + \frac{1}{\Gamma(3)}x^2 + \frac{1}{\Gamma(4)}x^3 - \frac{2}{\Gamma(5)}x^4 + \frac{1}{\Gamma(6)}x^5 + \frac{1}{\Gamma(7)}x^6 - \frac{2}{\Gamma(8)}x^7 \\
 &+ \frac{1}{\Gamma(9)}x^8 + \frac{1}{\Gamma(5/2)}x^{3/2} - \frac{2}{\Gamma(7/2)}x^{5/2} + \frac{1}{\Gamma(9/2)}x^{7/2} + \frac{1}{\Gamma(11/2)}x^{9/2} - \frac{2}{\Gamma(13/2)}x^{11/2} \\
 &+ \frac{1}{\Gamma(15/2)}x^{13/2} + \frac{1}{\Gamma(17/2)}x^{15/2} - \frac{2}{\Gamma(19/2)}x^{17/2} + \dots
 \end{aligned} \tag{8}$$

**5. Homotopy Analysis Method (HAM):**

Let us consider the fractional differential equation

$$\aleph(u(x)) = 0$$

Where  $\aleph$  is a nonlinear operator,  $x$  is independent variable,  $u(x)$  is an unknown function, respectively. For simplicity, we ignore all boundary or initial conditions, which can be treated in the similar way. By means of generalizing the traditional homotopy method, Liao (2004, 2012) <sup>[29, 31]</sup> constructs the so-called zero-order deformation equation

$$(1-q)L[\phi(x;q) - u_0(x)] = q\hbar\aleph(\phi(x;q)) \tag{9}$$

Here  $q \in [0,1]$  is the embedding parameter,  $\hbar \neq 0$  is a convergence-control parameter to be chosen later on,  $L$  is an auxiliary linear differential operator of integer order and possesses the property  $L(C) = 0$ , where  $C$  and  $\phi = \phi(x;q)$  are constant and unknown function respectively.

Expanding  $\phi(x;q)$  in Taylor series with respect to  $q$ , we have

$$\phi(x; q) = u_0(x) + \sum_{m=1}^{\infty} u_m(x) q^m \tag{10}$$

where 
$$u_m(x) = \frac{1}{m!} \left[ \frac{\partial^m \phi(x; q)}{\partial q^m} \right]_{q=0}$$

Differentiating the equation (9)  $m$ -times with respect to the embedding parameter  $q$  and then setting  $q = 0$ , and finally dividing them by  $m!$ , we have the  $m^{th}$ -order deformation equation

$$L[u_m(x) - \lambda_m u_{m-1}(x)] = \hbar R_m [\vec{u}_{m-1}(x)] \tag{11}$$

Where 
$$R_m [\vec{u}_{m-1}(x)] = \frac{1}{(m-1)!} \left[ \frac{\partial^{m-1} \mathfrak{N}(\phi(x; q))}{\partial q^{m-1}} \right]_{q=0}, \quad \lambda_m = \begin{cases} 0 & ; m \leq 1 \\ 1 & ; m \geq 2 \end{cases}$$

and 
$$\vec{u}_{m-1}(x) = [u_0(x), u_1(x), u_2(x), u_3(x), \dots, u_{m-1}(x)]$$

We define the equation (1), in terms of linear and nonlinear operators as follows:

$$L[\phi(x; q)] = D\phi(x; q) \tag{12}$$

and 
$$\mathfrak{N}[\phi(x; q)] = D\phi(x; q) + D_c^{1/2} \phi(x; q) + \phi(x; q) - x \tag{13}$$

where 
$$D = \frac{d}{dx}$$

Using (9), we construct the zeroth order deformation equation as

$$(1-q)L[\phi(x; q) - y_0(x)] = q\hbar \mathfrak{N}[\phi(x; q)] \tag{14}$$

Obviously, when  $q = 0$  and  $q = 1$  respectively, we can write from equation (14) the following relations:

$$\phi(x; 0) = y_0(x) = y(0) \quad \text{and} \quad \phi(x; 1) = y(x) \tag{15}$$

According to equations (10) and (11), we obtain the  $m^{th}$  order deformation equation as

$$L[y_m(x) - \lambda_m y_{m-1}(x)] = \hbar R_m [\vec{y}_{m-1}(x)] \tag{16}$$

where

$$\begin{aligned}
 R_m [\bar{y}_{m-1}(x)] &= \frac{1}{(m-1)!} \left[ \frac{\partial^{m-1} \mathfrak{R}[\phi(x; q)]}{\partial q^{m-1}} \right]_{q=0} \\
 &= \frac{1}{(m-1)!} \left[ \frac{\partial^{m-1}}{\partial q^{m-1}} \left\{ D\phi(x; q) + D_c^{1/2} \phi(x; q) + \phi(x; q) - x \right\} \right]_{q=0} \\
 &= Dy_{m-1}(x) + D_c^{1/2} y_{m-1}(x) + y_{m-1}(x) - x(1 - \lambda_m)
 \end{aligned} \tag{17}$$

From equations (16) and (17), we can write

$$L[y_m(x) - \lambda_m y_{m-1}(x)] = \hbar \left[ Dy_{m-1}(x) + D_c^{1/2} y_{m-1}(x) + y_{m-1}(x) - x(1 - \lambda_m) \right] \tag{18}$$

For  $m = 1$ , equation (18) becomes

$$L[y_1(x)] = \hbar \left[ Dy_0(x) + D_c^{1/2} y_0(x) + y_0(x) - x \right] \tag{19}$$

For  $m \geq 2$ , equation (18) reduces to

$$y_m(x) = y_{m-1}(x) + \hbar I \left[ Dy_{m-1}(x) + D_c^{1/2} y_{m-1}(x) + y_{m-1}(x) \right] \tag{20}$$

Taking  $\hbar = -1$ , we get from equations (1),(15), (19) and (20), after some rigorous calculations,

$$\begin{aligned}
 y_0(x) = y(0) = 1, \quad y_1(x) = -x + \frac{1}{2}x^2, \quad y_2(x) = \frac{1}{2}x^2 - \frac{1}{6}x^3 + \frac{1}{\Gamma\left(\frac{5}{2}\right)}x^{3/2} - \frac{1}{\Gamma\left(\frac{7}{2}\right)}x^{5/2}, \\
 y_3(x) = -\frac{1}{2}x^2 - \frac{1}{\Gamma\left(\frac{7}{2}\right)}x^{5/2} + \frac{1}{3\Gamma\left(\frac{9}{2}\right)}x^{7/2} + \frac{1}{24}x^4, \\
 y_4(x) = \frac{1}{\Gamma\left(\frac{7}{2}\right)}x^{5/2} + \frac{3}{\Gamma(4)}x^3 - \frac{2}{\Gamma(5)}x^4 + \frac{2}{\Gamma\left(\frac{9}{2}\right)}x^{7/2} - \frac{3}{\Gamma\left(\frac{11}{2}\right)}x^{9/2} - \frac{1}{\Gamma(6)}x^5, \\
 y_5(x) = -\frac{1}{\Gamma(4)}x^3 - \frac{5}{\Gamma(5)}x^4 - \frac{4}{\Gamma\left(\frac{9}{2}\right)}x^{7/2} + \frac{5}{\Gamma(6)}x^5 + \frac{4}{\Gamma\left(\frac{13}{2}\right)}x^{11/2} + \frac{1}{\Gamma(7)}x^6, \\
 y_6(x) = \frac{1}{\Gamma\left(\frac{9}{2}\right)}x^{7/2} + \frac{9}{\Gamma\left(\frac{11}{2}\right)}x^{9/2} + \frac{5}{\Gamma(5)}x^4 + \frac{5}{\Gamma(6)}x^5 - \frac{5}{\Gamma\left(\frac{13}{2}\right)}x^{11/2} \\
 - \frac{9}{\Gamma(7)}x^6 - \frac{5}{\Gamma\left(\frac{15}{2}\right)}x^{13/2} - \frac{1}{\Gamma(8)}x^7,
 \end{aligned}$$

$$y_7(x) = -\frac{1}{\Gamma(5)}x^4 - \frac{14}{\Gamma(6)}x^5 - \frac{6}{\Gamma(11/2)}x^{9/2} - \frac{14}{\Gamma(13/2)}x^{11/2} + \frac{14}{\Gamma(15/2)}x^{13/2} + \frac{14}{\Gamma(8)}x^7$$

$$+ \frac{6}{\Gamma(17/2)}x^{15/2} + \frac{1}{\Gamma(9)}x^8$$

and so on.

By using (10), we obtain

$$y(x) = y_0(x) + \sum_{m=1}^{\infty} y_m(x) \quad ; \text{ for } q=1 \tag{21}$$

Using the above values of  $y_n(x)$ ;  $n = 0, 1, 2, 3, \dots$  in (21), the solution of (1) is obtained as

$$y(x) = 1 - x + \frac{1}{2}x^2 + \frac{1}{6}x^3 - \frac{2}{24}x^4 - \frac{5}{\Gamma(6)}x^5 - \frac{8}{\Gamma(7)}x^6 + \frac{13}{\Gamma(8)}x^7 + \frac{1}{\Gamma(9)}x^8 + \frac{1}{\Gamma(5/2)}x^{3/2}$$

$$- \frac{1}{\Gamma(7/2)}x^{5/2} - \frac{2}{3\Gamma(9/2)}x^{7/2} - \frac{15}{\Gamma(13/2)}x^{11/2} + \frac{9}{\Gamma(15/2)}x^{13/2} + \frac{6}{\Gamma(17/2)}x^{15/2} + \dots \tag{22}$$

**6. Adomian Decomposition Method (ADM):**

Let us consider the differential equation in the following form (see Adomian, 1994, 1988; Wazwaz, 1999, 2001, 2002, 2005) [4, 3, 44, 47, 48, 49, 50].

$$L[y(x)] + N[y(x)] = f(x)$$

where  $L$  is a linear differential operator,  $N$  is a nonlinear operator,  $f(x)$  is a known function of  $x$ , respectively. Adomian decomposition method (ADM) defines the unknown function  $y(x)$  by an infinite series (see Adomian, 1994, 1988; Wazwaz, 1999, 2001, 2002, 2005) [4, 3, 44, 47, 48, 49, 50].

$$y(x) = \sum_{i=0}^{\infty} y_i(x) \tag{23}$$

where the components  $y_i(x)$  are determined by using recurrence relations.

The nonlinear part  $N[y(x)]$  can be decomposed into an infinite series of adomian polynomials given by (see Adomian, 1994, 1988; Wazwaz, 1999, 2001, 2002, 2005) [4, 3, 44, 47, 48, 49, 50]

$$N[y(x)] = \sum_{i=0}^{\infty} A_i(y_1, y_2, y_3, \dots, y_i) \tag{24}$$

where  $A_i$ 's are the so called adomian polynomials and defined by

$$A_i(y_1, y_2, y_3, \dots, y_i) = \frac{1}{i!} \left\{ \frac{d^i}{d\lambda^i} \left[ N \left( \sum_{j=0}^i \lambda^j y_j(x) \right) \right] \right\}_{\lambda=0}, \text{ for } i = 0, 1, 2, \dots \tag{25}$$

Taking  $L = \frac{d}{dx}$  and  $N[y(x)] = D_c^{1/2}y(x) + y(x)$ , let us rewrite the equation (1) as

$$L[y(x)] = -N[y(x)] + x \tag{26}$$

Applying  $L^{-1}$  on both sides of equation (26), we get

$$y(x) = \left( y(0) + \frac{1}{2}x^2 \right) - L^{-1} \{ N[y(x)] \} \tag{27}$$

Let us choose

$$y_0(x) = y(0) + \frac{1}{2}x^2 = 1 + \frac{1}{2}x^2 \tag{28}$$

Then after using equations (23) and (24) in (27), we get

$$y_{i+1}(x) = -L^{-1} \left[ A_i(y_0, y_1, y_2, \dots, y_i) \right], \text{ for } i = 0, 1, 2, \dots$$

where  $A_i$ 's are given by (25) and consequently, we obtain

$$y_1(x) = -\frac{x^3}{\Gamma(4)} - x - \frac{1}{\Gamma(7/2)} x^{5/2},$$

$$y_2(x) = \frac{1}{\Gamma(4)} x^3 + \frac{1}{\Gamma(5/2)} x^{3/2} + \frac{2}{\Gamma(9/2)} x^{7/2} + \frac{1}{\Gamma(3)} x^2 + \frac{1}{\Gamma(5)} x^4,$$

$$y_3(x) = -\left[ \frac{1}{\Gamma(9/2)} x^{7/2} + \frac{3}{\Gamma(5)} x^4 + \frac{2}{\Gamma(7/2)} x^{5/2} + \frac{3}{\Gamma(11/2)} x^{9/2} + \frac{1}{\Gamma(3)} x^2 + \frac{1}{\Gamma(4)} x^3 + \frac{1}{\Gamma(6)} x^5 \right],$$

$$y_4(x) = \frac{2}{\Gamma(5)} x^4 + \frac{3}{\Gamma(4)} x^3 + \frac{6}{\Gamma(6)} x^5 + \frac{1}{\Gamma(7)} x^6 + \frac{4}{\Gamma(11/2)} x^{9/2} + \frac{1}{\Gamma(7/2)} x^{5/2} + \frac{3}{\Gamma(9/2)} x^{7/2} + \frac{4}{\Gamma(13/2)} x^{11/2},$$

$$y_5(x) = - \left[ \frac{1}{\Gamma(4)} x^3 + \frac{6}{\Gamma(5)} x^4 + \frac{6}{\Gamma(6)} x^5 + \frac{10}{\Gamma(7)} x^6 + \frac{1}{\Gamma(8)} x^7 + \frac{4}{\Gamma(9/2)} x^{7/2} + \frac{5}{\Gamma(11/2)} x^{9/2} + \frac{10}{\Gamma(13/2)} x^{11/2} + \frac{5}{\Gamma(15/2)} x^{13/2} \right],$$

$$y_6(x) = \frac{5}{\Gamma(5)} x^4 + \frac{11}{\Gamma(6)} x^5 + \frac{16}{\Gamma(7)} x^6 + \frac{15}{\Gamma(8)} x^7 + \frac{1}{\Gamma(9)} x^8 + \frac{1}{\Gamma(9/2)} x^{7/2} + \frac{10}{\Gamma(11/2)} x^{9/2} + \frac{11}{\Gamma(13/2)} x^{11/2} + \frac{20}{\Gamma(15/2)} x^{13/2} + \frac{6}{\Gamma(17/2)} x^{15/2},$$

$$y_7(x) = - \left[ \frac{1}{\Gamma(5)} x^4 + \frac{15}{\Gamma(6)} x^5 + \frac{22}{\Gamma(7)} x^6 + \frac{36}{\Gamma(8)} x^7 + \frac{21}{\Gamma(9)} x^8 + \frac{1}{\Gamma(10)} x^9 + \frac{6}{\Gamma(11/2)} x^{9/2} + \frac{21}{\Gamma(13/2)} x^{11/2} + \frac{27}{\Gamma(15/2)} x^{13/2} + \frac{35}{\Gamma(17/2)} x^{15/2} + \frac{7}{\Gamma(19/2)} x^{17/2} \right]$$

and so on

By using (23) and above values of  $y_n(x)$ ;  $n = 0, 1, 2, \dots$  the solution of (1) is obtained as

$$y(x) = 1 - x + \frac{1}{2} x^2 + \frac{1}{\Gamma(4)} x^3 - \frac{2}{\Gamma(5)} x^4 - \frac{5}{\Gamma(6)} x^5 - \frac{15}{\Gamma(7)} x^6 - \frac{22}{\Gamma(8)} x^7 - \frac{20}{\Gamma(9)} x^8 - \frac{1}{\Gamma(10)} x^9 + \frac{1}{\Gamma(5/2)} x^{3/2} - \frac{2}{\Gamma(7/2)} x^{5/2} + \frac{1}{\Gamma(9/2)} x^{7/2} - \frac{16}{\Gamma(13/2)} x^{11/2} - \frac{12}{\Gamma(15/2)} x^{13/2} - \frac{29}{\Gamma(17/2)} x^{15/2} - \frac{7}{\Gamma(19/2)} x^{17/2} + \dots \tag{29}$$

**7. Solution of equation (1) by the Power Series Method (PSM):**

Let 
$$y(x) = \sum_{n=0}^{\infty} a_n x^{n/2} \tag{30}$$

be the power series solution of the fractional linear ordinary differential equation (1).

Applying  $I^{1/2}$  on both sides of (1), we obtain

$$I^{1/2} (Dy(x)) + I^{1/2} (D_c^{1/2} y(x)) + I^{1/2} (y(x)) = I^{1/2} (x) \tag{31}$$

Now substituting (30) in (31), one gets

$$\sum_{n=1}^{\infty} \frac{\Gamma\left(\frac{n}{2}+1\right)}{\Gamma\left(\frac{n}{2}+\frac{1}{2}\right)} a_n x^{\frac{n-1}{2}} + \sum_{n=1}^{\infty} a_n x^{\frac{n}{2}} + \sum_{n=0}^{\infty} \frac{\Gamma\left(\frac{n}{2}+1\right)}{\Gamma\left(\frac{n}{2}+\frac{3}{2}\right)} a_n x^{\frac{n+1}{2}} = \frac{x^{\frac{3}{2}}}{\Gamma\left(\frac{5}{2}\right)} \quad (32)$$

With  $a_0 = 1$ , equating the co-efficient of different powers of  $x$ , we obtain

$$a_1 = 0, a_2 = -1, a_3 = \frac{1}{\Gamma\left(\frac{5}{2}\right)}, a_4 = \frac{1}{2}, a_5 = -\frac{2}{\Gamma\left(\frac{7}{2}\right)}, a_6 = \frac{1}{6}, a_7 = \frac{1}{\Gamma\left(\frac{9}{2}\right)}, a_8 = -\frac{2}{\Gamma(5)},$$

$$a_9 = \frac{1}{\Gamma\left(\frac{11}{2}\right)}, a_{10} = \frac{1}{\Gamma(6)}$$

and so on.

Then the Power Series Solution of (1) is

$$y(x) = 1 - x + \frac{1}{2}x^2 + \frac{1}{\Gamma(4)}x^3 - \frac{2}{\Gamma(5)}x^4 + \frac{1}{\Gamma(6)}x^5 + \frac{1}{\Gamma\left(\frac{5}{2}\right)}x^{\frac{3}{2}} - \frac{2}{\Gamma\left(\frac{7}{2}\right)}x^{\frac{5}{2}}$$

$$+ \frac{1}{\Gamma\left(\frac{9}{2}\right)}x^{\frac{7}{2}} + \frac{1}{\Gamma\left(\frac{11}{2}\right)}x^{\frac{9}{2}} + \dots \quad (33)$$

## 8. Conclusion

In this paper, four different methods- generalized differential transform method (GDMT), homotopy analysis method (HAM), adomian decomposition method (ADM) and power series method (PSM) applied to find the approximate solution of a linear fractional differential equation. It has been found that the solution obtained by these methods is significantly identical. The advantages of these global methodologies – GDTM, HAM and ADM lie in the fact that they not only lead to an analytical continuous approximation which is very rapidly convergent, but also show the dependence, giving insight into the character and behavior of the solution just as in a closed form solution. Furthermore, these methods do not require any transformation technique, linearization or discretization of the variables and they do not make closure approximation or smallness assumption. These techniques may also be applied to the solution of the nonlinear ordinary differential equations, nonlinear partial differential equations and nonlinear integral equations.

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